

Noise-driven adaptation: *in vitro* and
mathematical analysis

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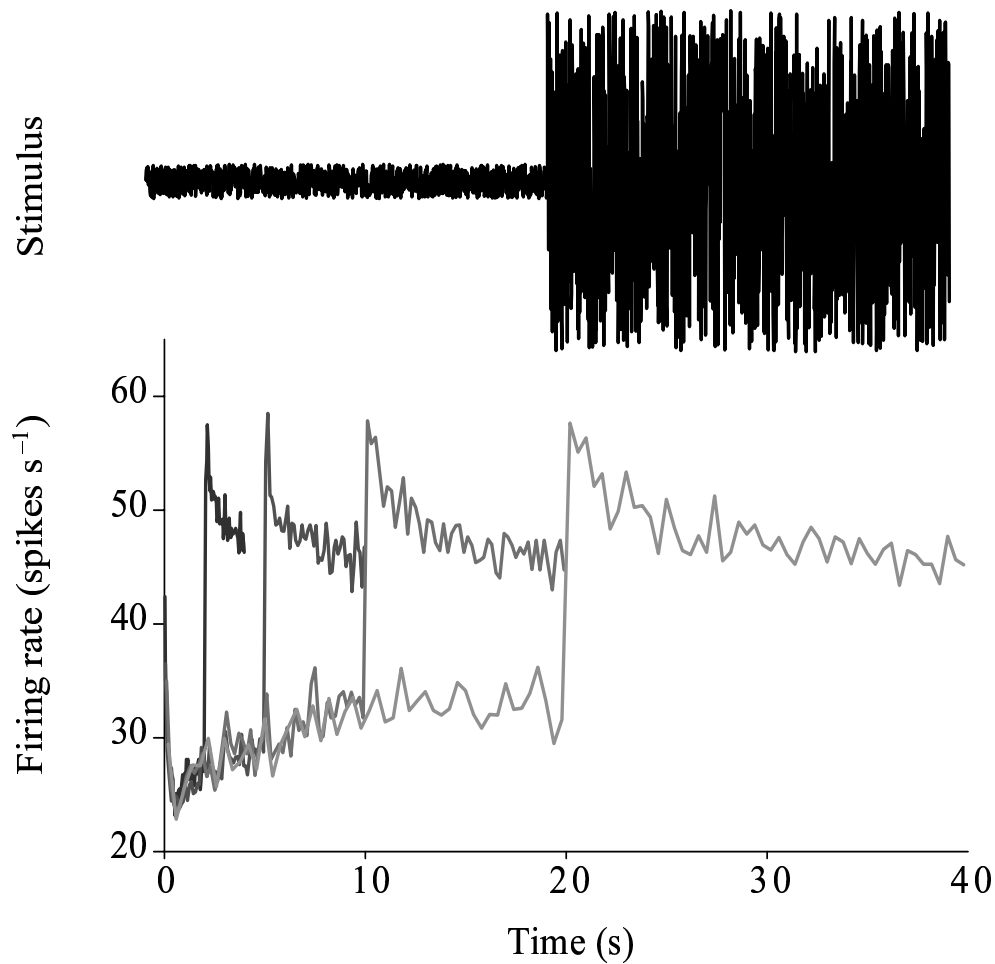
Introduction

Observation: changes in scale (not just mean) of random input can lead to “adaptation” (defined in more depth below) [13, 12, 5]

2 hypotheses:

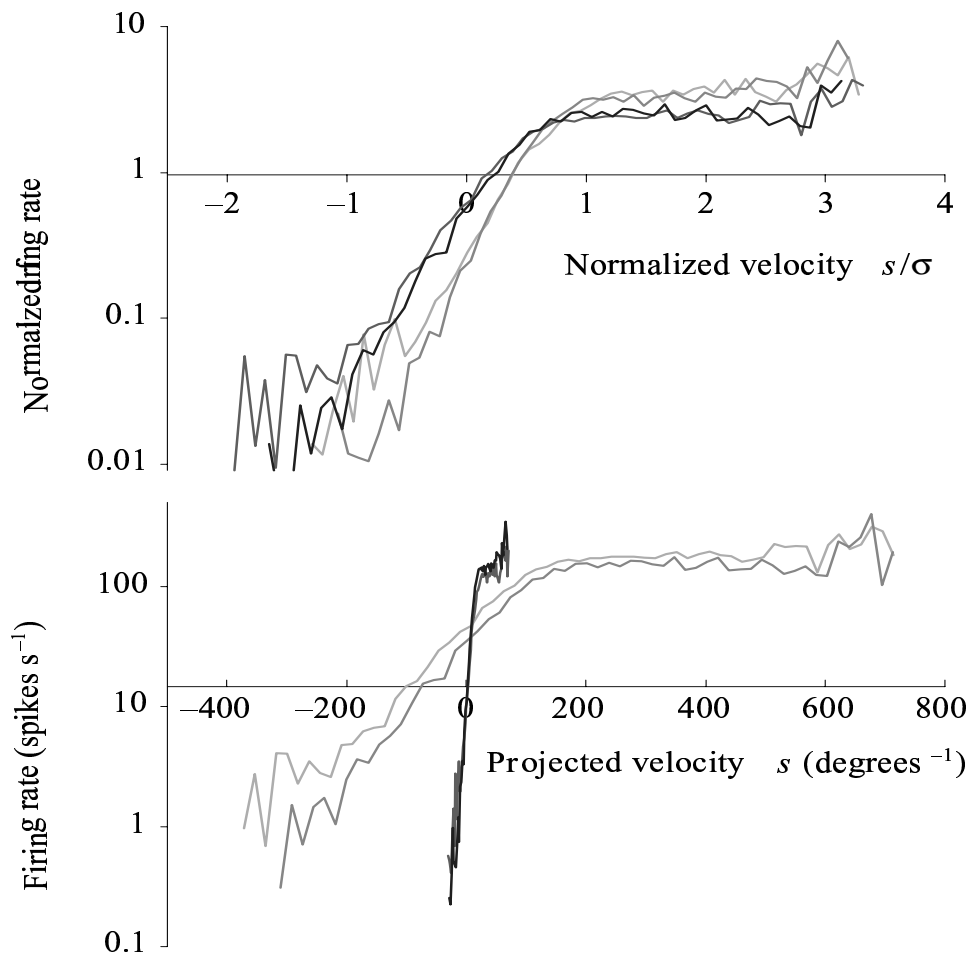
- (physiology) adaptation is at least partially single-cell phenomenon
- (math) adaptation is generic (independent of cellular details)

Results of Fairhall et al. (1)



- firing rate “adapts” to changes in input scale
- adaptation to upward jumps is faster than to downward jumps
- timescale of adaptation depends on timescale of scale changes

Results of Fairhall et al. (2)



- N -function adapts (exact definition below)

Experimental methods

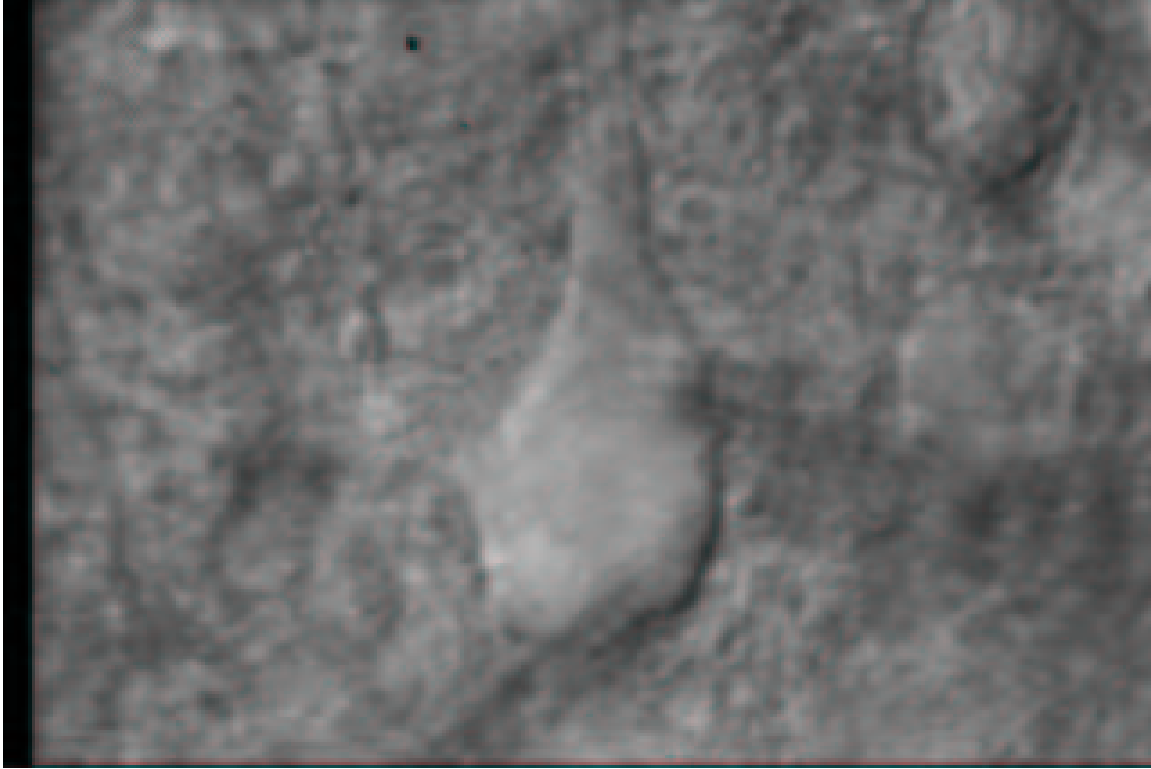
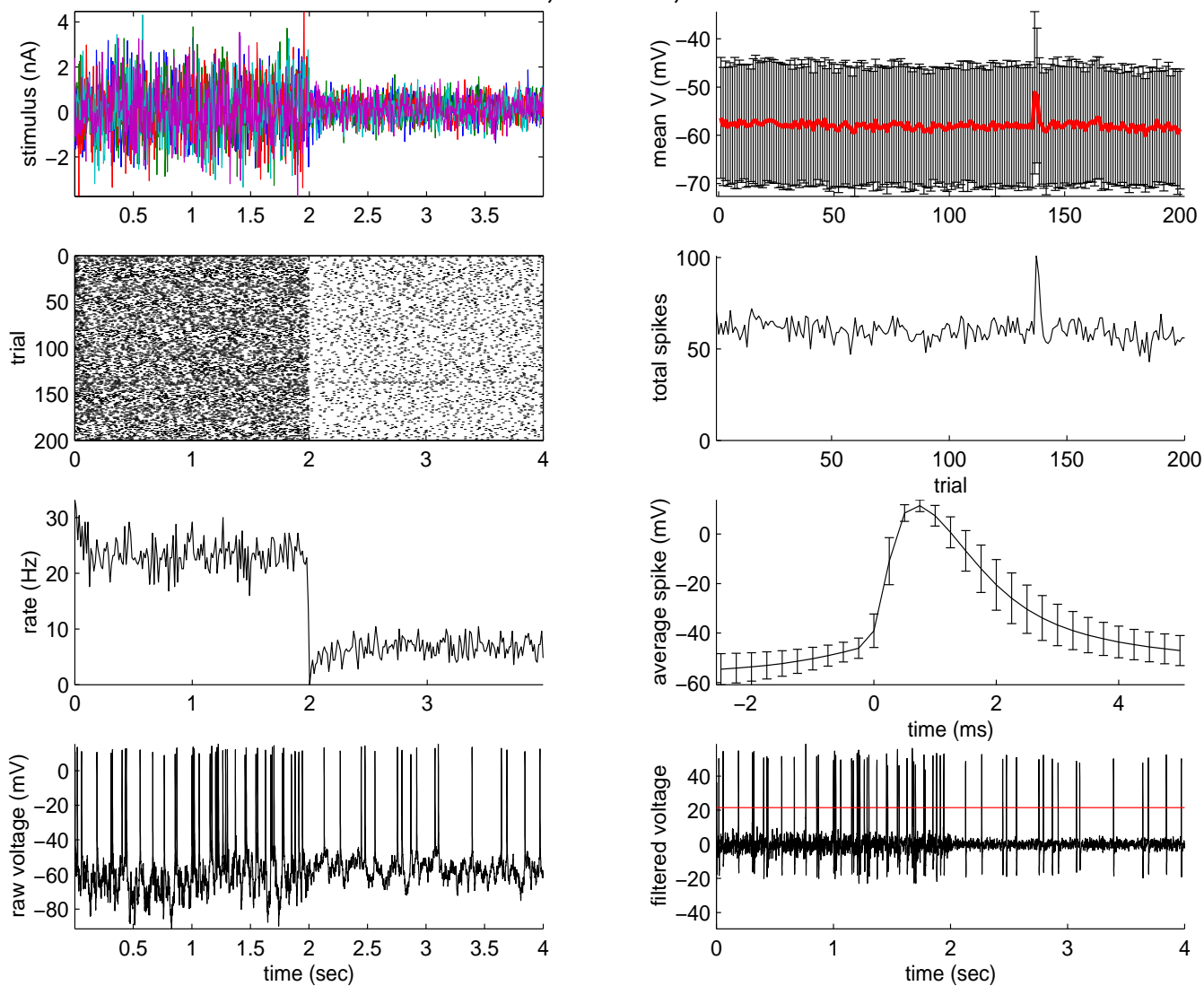


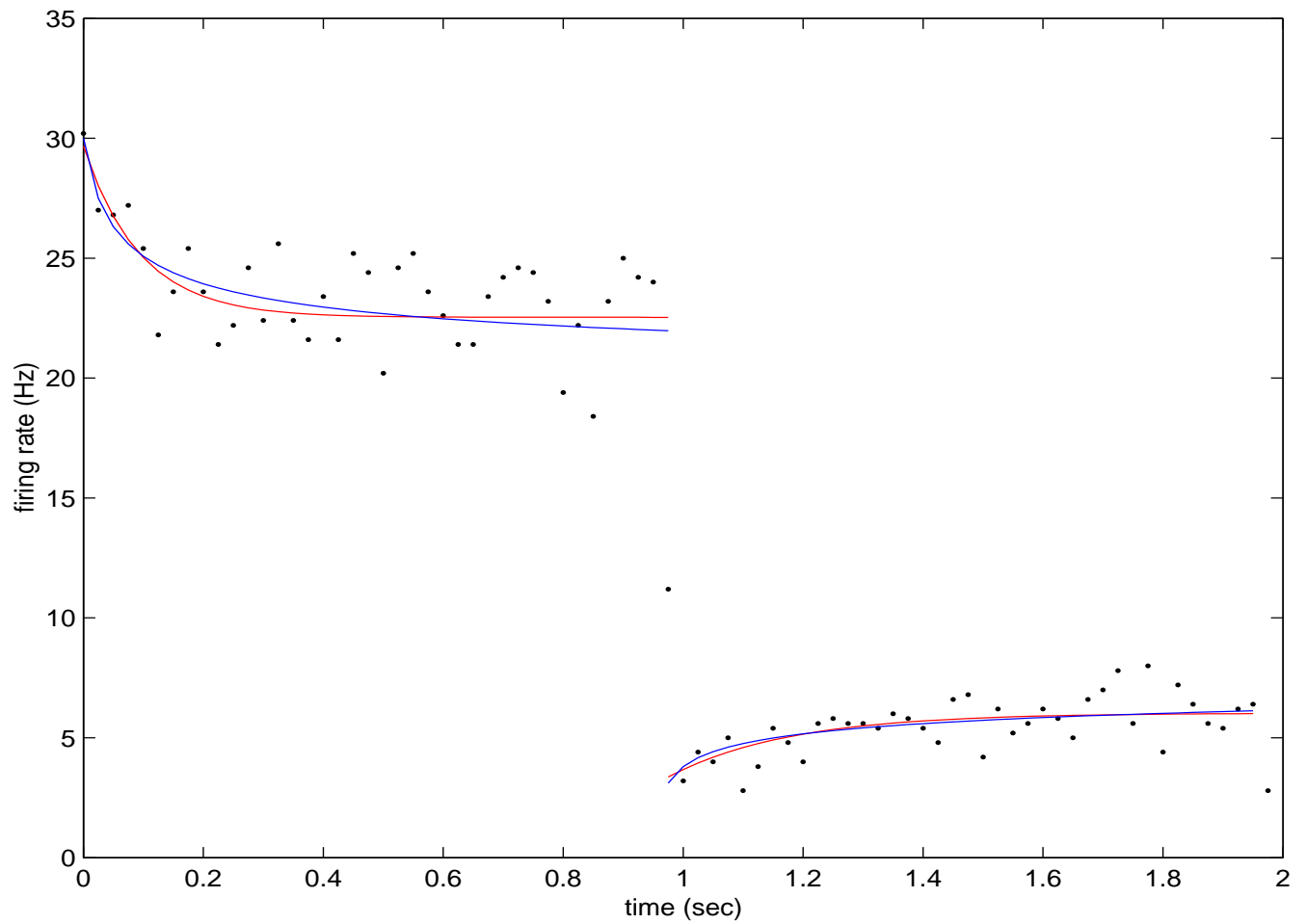
Figure 1: Sagittal slices were prepared from adolescent and adult rats (P14-P24) as described in [10]. Briefly, slices were maintained at 30°C in artificial cerebrospinal fluid consisting of (in mM): 125 NaCl, 2.5 KCl, 25 glucose, 25 NaHCO₃, 1.25 NaH₂PO₄, 2 CaCl₂, and 1 MgCl₂. Cells were visualized using infrared differential interference contrast microscopy with a 40X water immersion objective. Dual-electrode whole-cell recordings were made using pipettes with 5-15 M resistance when filled with (in mM): 100 K-gluconate, 20 KCl, 4 ATP-Mg, 10 phosphocreatine, 0.3 GTP, and 10 HEPES, pH 7.3 (310 mOsm). Recordings were performed in current clamp using Axoclamp 2B amplifiers (Axon Instruments, Foster City, CA), and stimulus presentation and data acquisition was managed using IGOR (Wavemetrics, Lake Oswego, OR). Gaussian noise current stimuli were delivered through one electrode, while voltage was recorded through the other electrode and processed on- and off-line (note that the noise current was not “frozen,” that is, a new noise current was drawn i.i.d. for each trial). Panel shows a photograph of a cell with the recording and stimulating electrodes partially visible.

Basic experimental data

cell 84, run 20, SQN

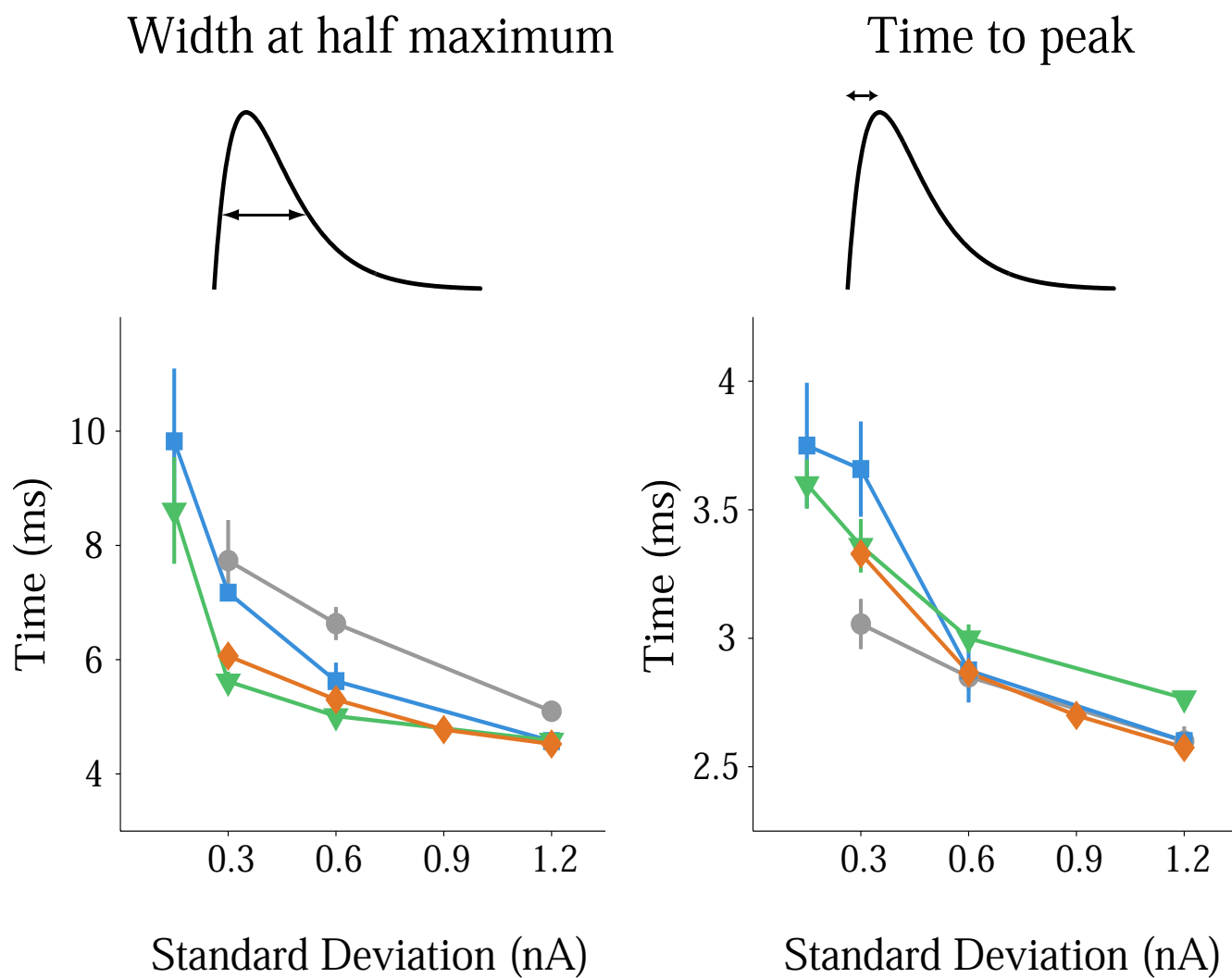


Experimental Result 1: rate adaptation



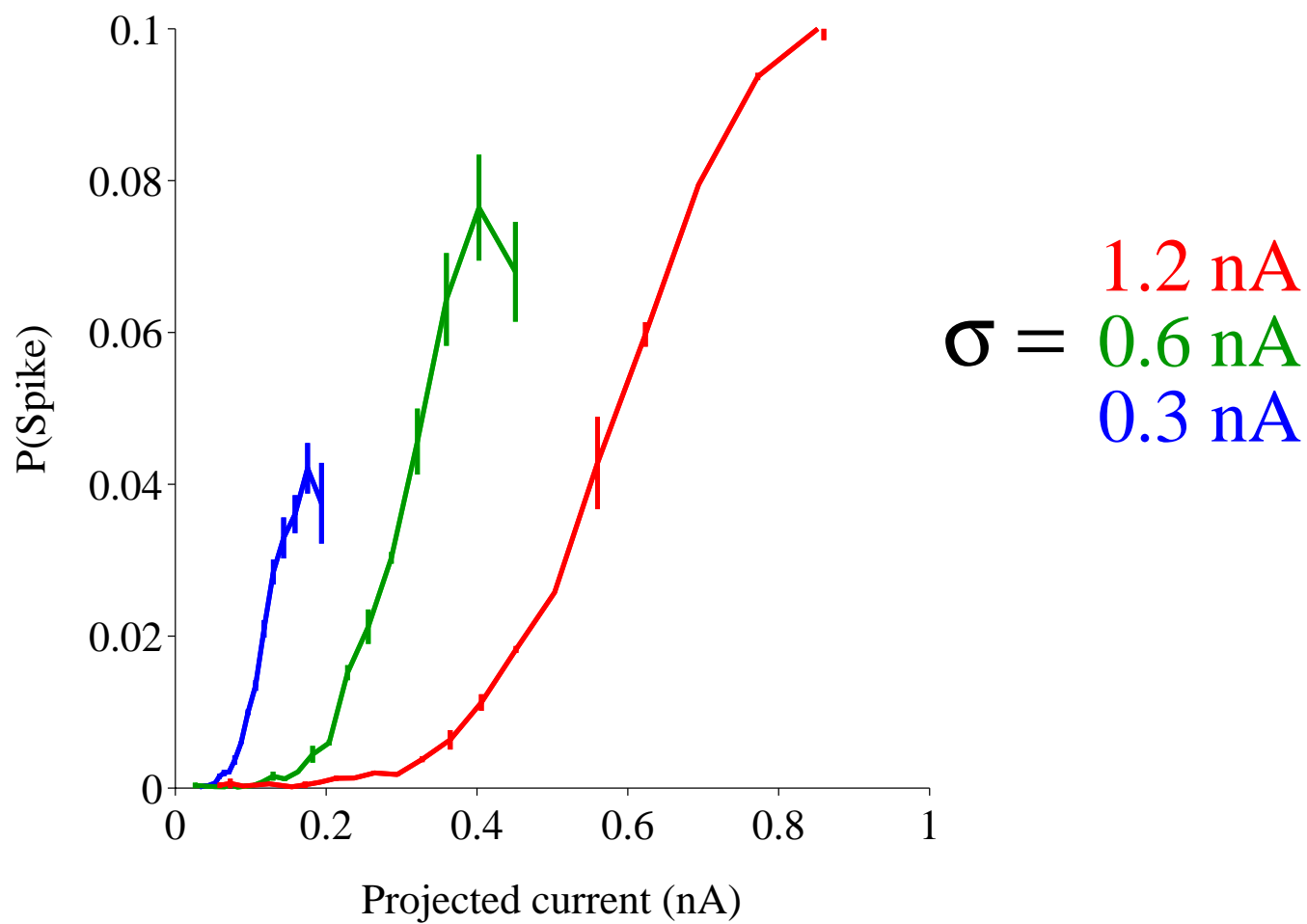
- $\tau_- > \tau_+$
- No oscillations seen
- No period-dependence seen (c.f. Fairhall et al.)

Result 2: changes in spike-triggered average



C.f. [3, 1, 13]

Result 3: adaption of N -function



Definition of N -function

$$N(K\vec{x}) \equiv P(\textit{spike}|K\vec{x})$$

(\vec{x} is stimulus, K is projection onto cell's “linear filter,” estimated by spike-triggered average)

$K\vec{x}$ = “projected current”

Theoretical methods: IF model

$$\frac{dV}{dt} = \frac{1}{\tau_m}(V_L - V + R_m I) - (V_{th} - V_{reset})\delta(V - V_{th})$$

τ_m = membrane time constant ≈ 20 ms

V_L = leak reversal potential ≈ -70 mV

R_m = membrane resistance $\approx 40 M\Omega$

I = input current (white Gaussian noise, fixed DC)

V_{th} = threshold potential ≈ -55 mV

V_{reset} = reset potential ≈ -65 mV

Basic tool: Fokker-Planck equation

Basic idea: instead of modeling single cells, model population [7, 2, 9]

Fundamental object: $P(V)$, probability distribution on voltage

$$\frac{\partial P}{\partial t} = L(P) + F(t)(\delta(V - V_{reset}) - \delta(V - V_{th})),$$

$$L \equiv \frac{\sigma_0^2}{2} \frac{\partial^2 P}{\partial V^2} + \frac{1}{\tau_m} \frac{\partial[(V - V_0)P]}{\partial V};$$

$F(t)$ time-dependent mean firing rate of the cell

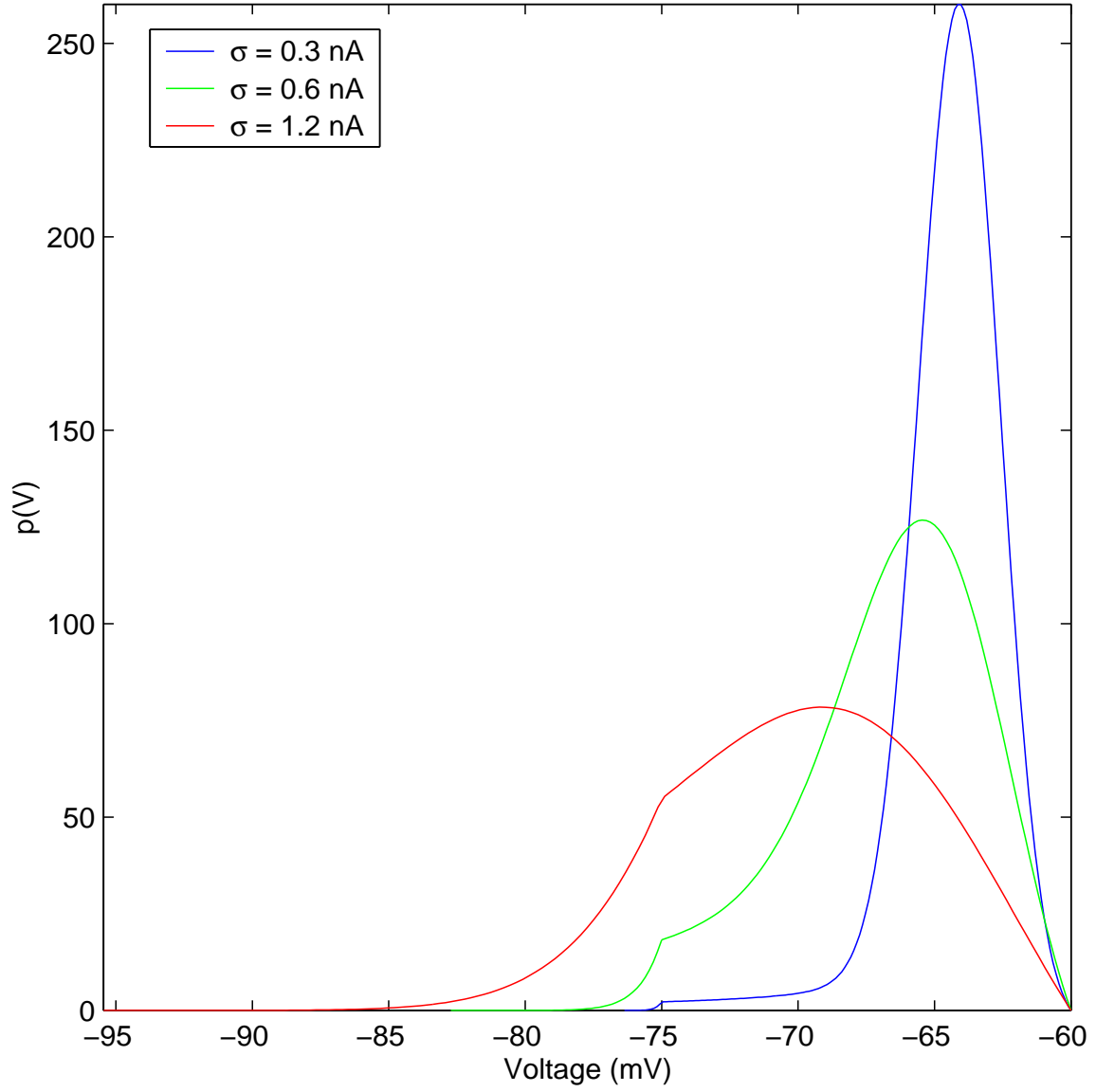
$$\sigma_0 \equiv R_m \sigma / \tau_m$$

$$V_0 \equiv R_m \mu + V_L$$

This equation exact in IF framework — no approximations made

Steady-state solution to FP equation [8, 2]

$$P(V) = \frac{2F}{\sigma_0^2} \int_{\max(V, V_{reset})}^{V_{th}} dV' e^{\frac{(V'-V_0)^2 - (V-V_0)^2}{\tau_m \sigma_0^2}}$$



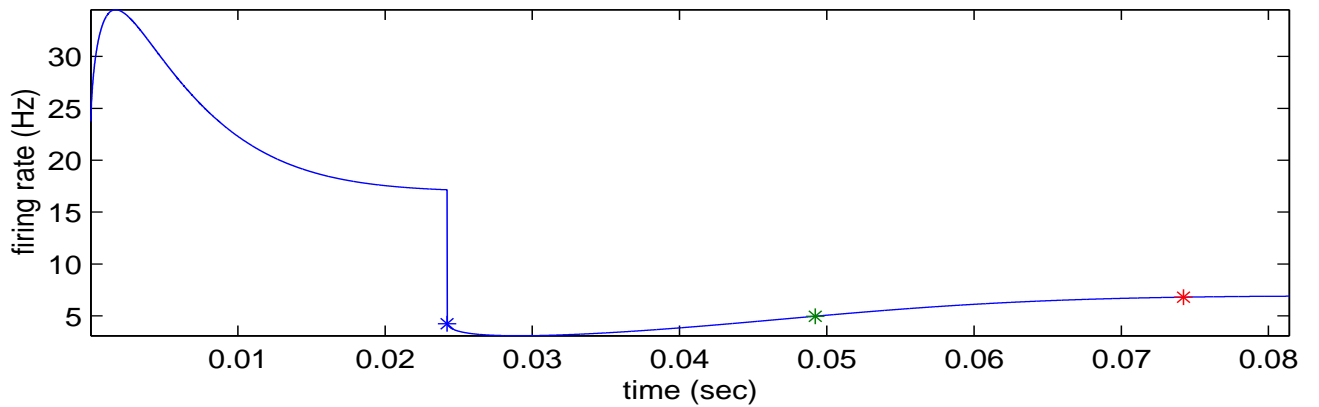
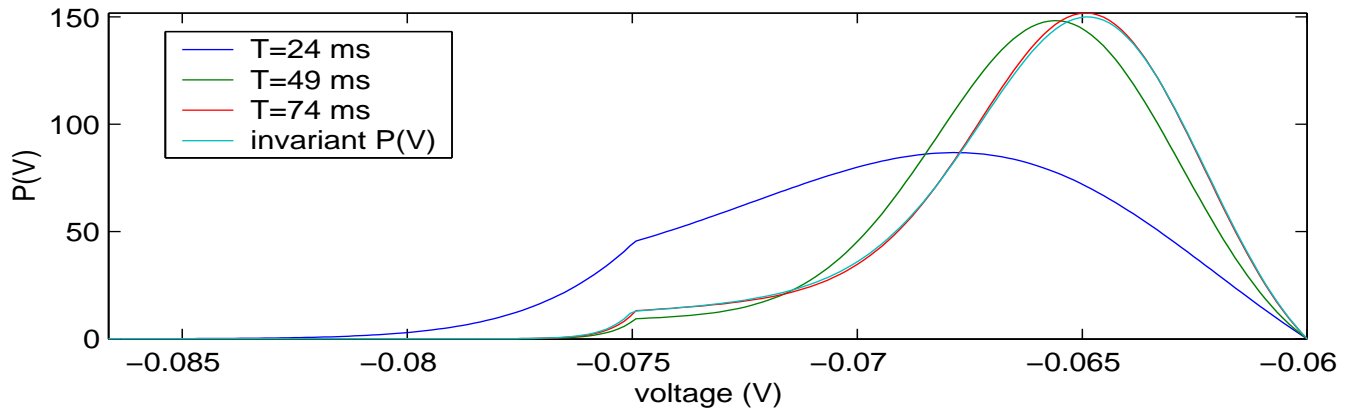
Theoretical result 1: rate adaptation

$$F_{\infty} = -\frac{\sigma_0^2}{2} \frac{\partial P_{\infty}(V)}{\partial V}$$

$$F_0 = -\frac{\sigma_0^2}{2} \frac{\partial P_0(V)}{\partial V}$$

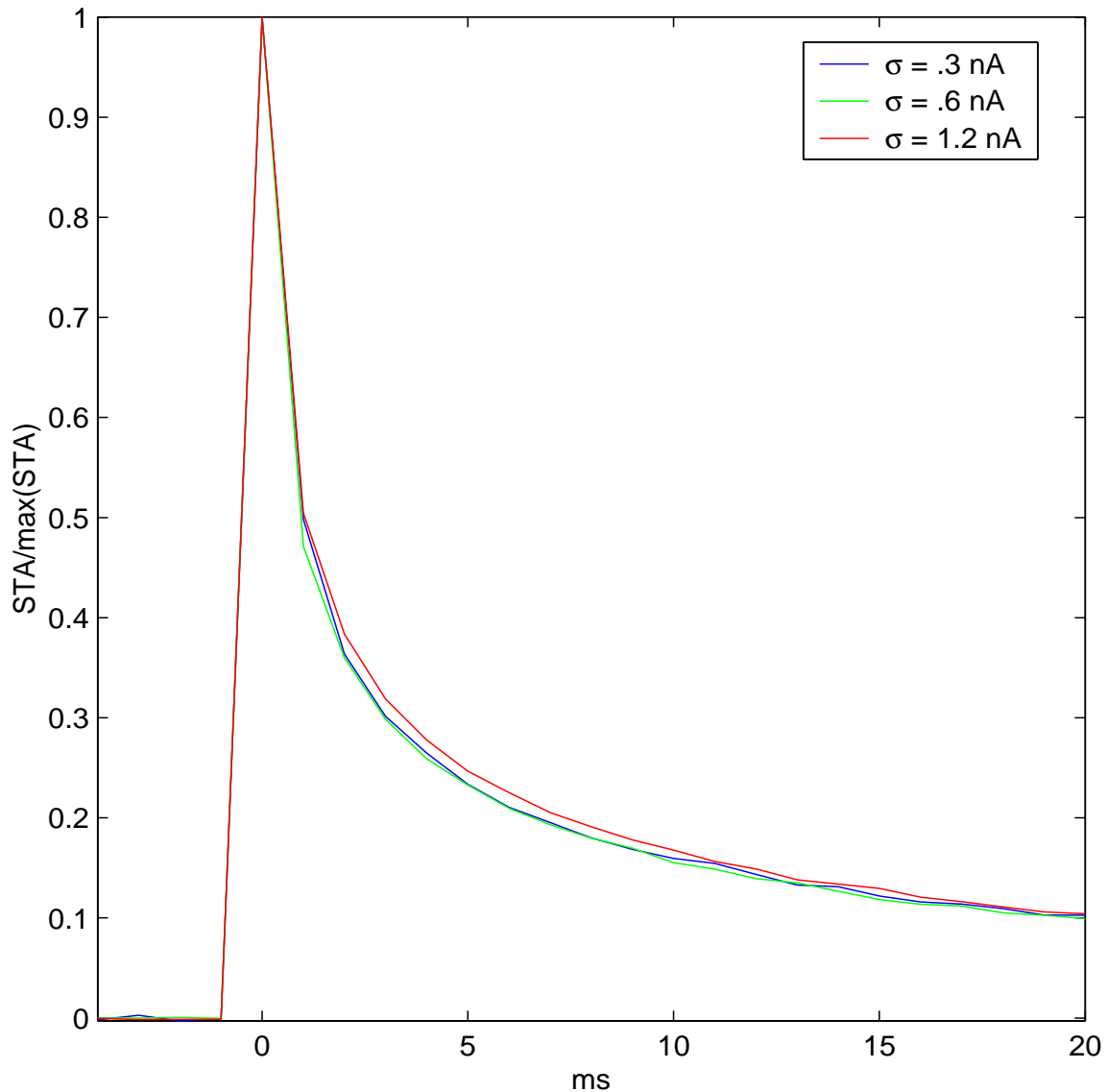
$$\frac{\partial^i F_0}{\partial t^i} = -\frac{\sigma_0^2}{2} \frac{\partial (L^i(P_0))}{\partial V}$$

(but F is not analytic; power series expansion invalid)



Result 2: spike-triggered average

Does not show width scaling behavior seen *in vitro*; different model necessary (see, e.g., [6])



Some exact theory possible, using invariant measure of FP equation; STA scales with σ and (more surprisingly) time step of numerical integration

Result 3: adaptation of three N -functions

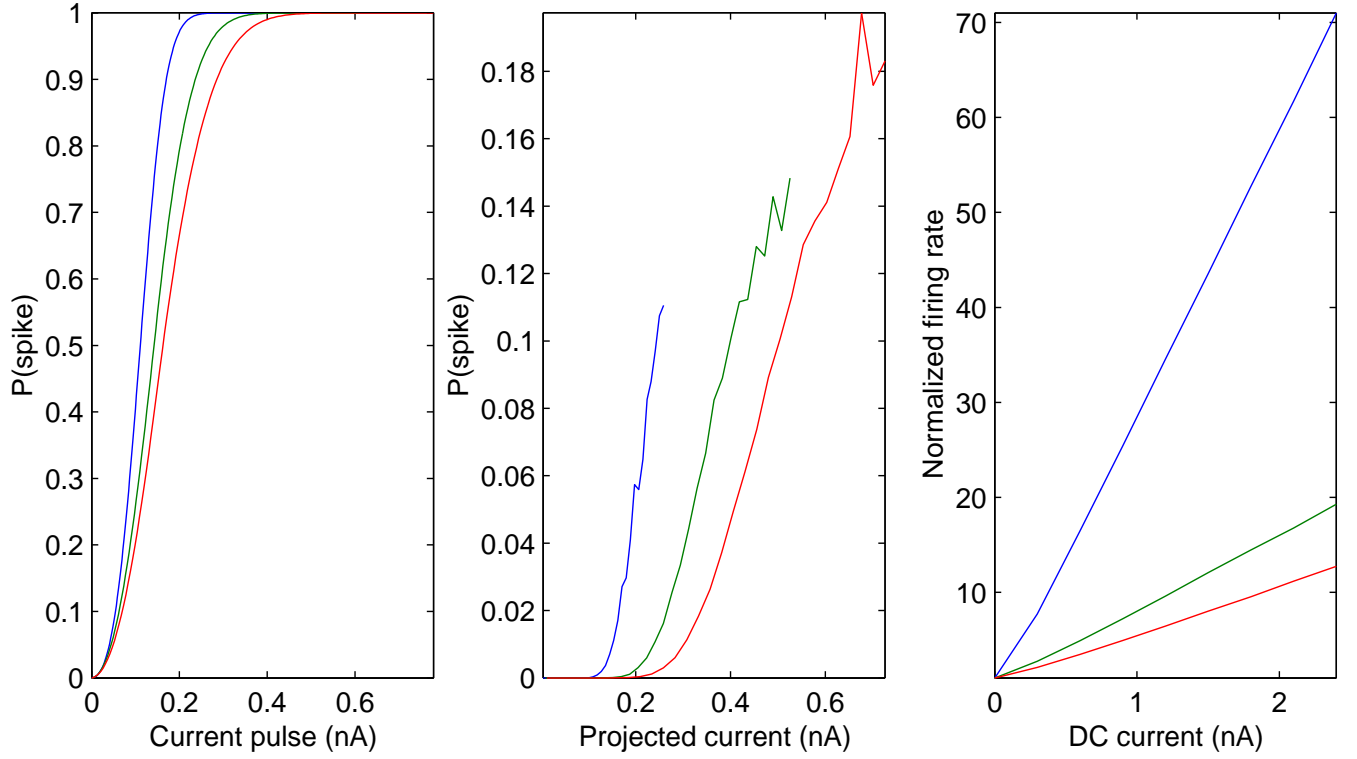


Figure 2: Three “gain functions” for integrate-and-fire cell: middle panel shows N -function, computed by Monte Carlo; left panel is “transient” function $F_0(x, \sigma)$, and right is “long-time” function $F_\infty(x, \sigma)/F_\infty(0, \sigma)$, both computed analytically. (Note that $F_\infty(x, \sigma)/F_\infty(0, \sigma)$ is normalized so that the y-axis is a dimensionless ratio.)

IF gain functions: definitions

Transient gain function:

$$\begin{aligned} F_0(x, \sigma) &\equiv \lim_{dt \rightarrow 0} \lim_{T \rightarrow 0} P\left(\text{spike} \in (-dt, 0] \mid \int_{-T}^0 I(t) dt = x\right) \\ &= \int_{V_{th} - \frac{x R_m}{\tau_m}}^{V_{th}} P(V) dV, \end{aligned}$$

Long-time gain function:

$$\begin{aligned} F_\infty(x, \sigma) &\equiv \lim_{T \rightarrow \infty} P\left(\text{spike} \in (-dt, 0] \mid \int_{-T}^0 I(t) dt = xT\right) \\ &\approx -\frac{\sigma_0^2}{2} \frac{\partial P_{\mu+x}(V)}{\partial V} \Big|_{V=V_{th}} dt, \end{aligned}$$

x corresponds to the projected current (x-axis, Fig. 2b).

Conclusions

- Cortical somata display much of the adaptive behavior observed *in vivo*, in various species and preparations
- Much of this behavior is replicated in the simplest possible model [11]
- Different view of adaptive behavior; no efficient coding concepts invoked

Directions

- Dynamics of current-based FP equation
- Conductance-based FP equation (no second-order differential term; instead, jumps, like a delay equation in space).
- Dependence of τ on period in Fairhall et al. data: network effect?

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